

How to quantize the Dirac Field

$$\mathcal{L} = \bar{\Psi}(i\not{\partial} - m)\Psi = \Psi^\dagger \rho(i\vec{\partial} \cdot \vec{\gamma} + i\partial_0 \gamma^0 - m)\Psi$$

$$\frac{\partial \mathcal{L}}{\partial(\partial_0 \Psi)} = i\Psi^\dagger = \Pi_\Psi$$

$$\frac{\partial \mathcal{L}}{\partial(\partial_0 \Psi^\dagger)} = 0$$

For this form of the Lagrangian

$$H = \int d^3x (\Pi_\Psi \partial_0 \Psi - \mathcal{L}) = - \int d^3x \underbrace{\Psi^\dagger}_{\bar{\Psi}} \underbrace{\gamma^0 (i\vec{\partial} \cdot \vec{\gamma} - m)}_{i\not{\partial} - m} \Psi$$

If we let $\vec{\alpha} = \gamma^0 \vec{\gamma}$ and $\beta = \gamma^0$ then $H = \int d^3x \Psi^\dagger h_{\text{Dirac}} \Psi$

where $h_{\text{Dirac}} = -i\vec{\alpha} \cdot \vec{\nabla} + m\beta$

How to quantize?

Try $[\Psi_a(x), \underbrace{\Pi_b(y)}_{i\Psi^\dagger}]_{x^0=y^0} = i\delta^3(\vec{x}-\vec{y})\delta_b^a$ exactly the same as scalar field

$\left. \begin{matrix} a=1,2,3,4 \\ b=1,2,3,4 \end{matrix} \right\} 4 \text{ component objects}$

However, this introduces inconsistencies - it turns out that you need to use anti commutation relations

$$\Psi_a(\vec{x}) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} e^{i\vec{p} \cdot \vec{x}} \sum \left(\underset{\substack{\uparrow \\ \text{positive energy} \\ \text{soln}}}{a} a_{\vec{p}} u_a^s(\vec{p}) + b_{-\vec{p}}^s \underset{\substack{\uparrow \\ \text{negative energy soln}}}{V_a^s(-\vec{p})} \right)$$

Expand in terms of eigenfunctions of Dirac eqn. we want to impose a comm relation between a + b .

$$[a_{\vec{p}}^r, a_{\vec{q}}^{s\dagger}] = [b_{\vec{p}}^r, b_{\vec{q}}^s] = (2\pi)^3 \delta(\vec{p}-\vec{q}) \delta_{rs} \leftarrow \text{this is what we'd expect.}$$

But if we take $\Psi(\vec{x})$ and Π and try to compute the other commutation $[\Psi_a, \Pi_0]$

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$$+ [b_{-\bar{p}}^r, b_{-\bar{q}}^{s\dagger}] V_a^r(\bar{p}) \overline{V}_b^s(-\bar{q}) \Big) \gamma_{b'b}^0$$

assume $[a, b] = 0$. Plug in assumptions for commutativity

$$= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2E_p} e^{i\vec{p} \cdot (\vec{x} - \vec{y})} \sum_{r=0}^1 \left(U^r_{(\vec{p})} \bar{U}^r_{(\vec{p})} + V^r_{(-\vec{p})} \bar{V}^r_{(-\vec{p})} \right) \gamma^0$$

We wrote down identities for this kind of sum:

$$\left(\sum_r u(\vec{p}) \bar{u}(\vec{p}) \right)_{ab} = (\cancel{p} + m)_{ab} = (E_p \gamma^0 - \vec{p} \cdot \vec{\gamma} + m)_{ab}$$

$$\left(\sum_r \hat{V}(\vec{r}) \bar{V}^r(\vec{r}) \right)_{ab'} = (\not{\vec{p}} - m)_{ab'} = (E \gamma^0 + \vec{p} \cdot \vec{\gamma} - m)_{ab'}$$

add: $(\gamma^\mu \epsilon_\mu \gamma^\nu)$

$$[\psi(\bar{x}), \psi^\dagger(\bar{y})] = 2E_p \int \frac{d^3p}{(2\pi)^3} \frac{1}{2E_p} e^{i\vec{p} \cdot (\vec{x} - \vec{y})} 2E_p \underbrace{(\delta_{ab})^2}_{\delta_{ab}}$$

$$= \frac{\int d^3x}{\int (2\pi)^3} \delta^3(\bar{x} - \bar{y}) \mathbb{1}^{ab}$$

Compute hamiltonian

$$H = \int d^3x \psi^\dagger h_{\text{Dirac}} \psi \quad \begin{matrix} \text{lots of} \\ \text{work} \end{matrix} \quad \cancel{\int d^3p} \int \frac{d^3p}{(2\pi)^3} \sum_s E_p (a_p^{st} a_p^s - b_p^{st} b_p^s)$$

looks alot like spin 0 result but the minus sign is a problem

$$a_p |0\rangle = 0 \quad \cancel{b_p} a_p^\dagger |0\rangle = |\bar{p}\rangle$$

$$b_p |0\rangle = 0 \quad b_p^\dagger |0\rangle = |\bar{q}\rangle \quad \text{negative energy}$$

So $H|q\rangle = -E_q |q\rangle \rightarrow$ we can get arbitrarily negative energies
 so this is bad.

However, causality is good:

$$e^{iHt} a_p^s e^{-iHt} = a_p^s e^{-iE_p t}$$

$$b_p^s = b_p^s e^{+iE_p t}$$

Plug into $\psi(\bar{x})$:

$$\psi(\bar{x}) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s (a_p^s V_p^s(\bar{p}) e^{-ip \cdot \bar{x}} + b_{\bar{p}}^s V_p^s(\bar{p}) e^{+ip \cdot \bar{x}})$$

$$\bar{\psi}(\bar{y}) = \text{similar}$$

$$[\psi_a(\bar{x}), \bar{\psi}_b(\bar{y})] = (i\partial_x + m)_{ab} \underbrace{(D(\bar{x}-\bar{y}) - D(\bar{y}-\bar{x}))}_{[\phi(\bar{x}), \phi(\bar{y})]}$$

$\uparrow$
 commutes if space like separated

We got 2nd quantization + causality, but the energy is troublesome. You could fool around but there will always be a problem with one of these 3 things.

$$\langle 0 | [\psi_a, \bar{\psi}_b] | 0 \rangle = \langle 0 | \psi_a \bar{\psi}_b | 0 \rangle \leftarrow \text{causality works but } \psi \text{ is weird.}$$

The solution How to quantize Dirac field.

$$\psi(x) = \int \frac{d^3x}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_s (a_p^s U^s(\vec{p}) e^{-ip \cdot x} + b_{\vec{p}}^{\dagger} V^s(\vec{p}) e^{+ip \cdot x})$$

$$\bar{\psi} = \sum_s (b_{\vec{p}}^s \bar{V}^s(\vec{p}) e^{-ip \cdot x} + a_{\vec{p}}^{\dagger} \bar{U}^s(\vec{p}) e^{ip \cdot x})$$

Suppose $\{\psi(x), \pi_{\psi}(\bar{y})\}_{x=y} = i \delta^3(\vec{x}-\vec{y}) \delta$

S.:

$$H = \int \frac{d^3p}{(2\pi)^3} \sum_s E_p (a_{\vec{p}}^{\dagger} a_{\vec{p}}^s - b_{\vec{p}}^s b_{\vec{p}}^{\dagger})$$

$$= \int \frac{d^3p}{(2\pi)^3} \sum_s E_p (a_{\vec{p}}^{\dagger} a_{\vec{p}}^s + b_{\vec{p}}^{\dagger} b_{\vec{p}}^s - \{b_{\vec{p}}^s, b_{\vec{p}}^{\dagger}\})$$

~~Normal ordering:~~

$$b_{\vec{p}}^{\dagger} b_{\vec{p}}^s b_{\vec{q}}^r |0\rangle = b_{\vec{p}}^{\dagger} \{b_{\vec{p}}^s, b_{\vec{q}}^r\} |0\rangle = b_{\vec{p}}^{\dagger} b_{\vec{q}}^r b_{\vec{p}}^s |0\rangle$$

$$= (2\pi)^3 \delta^3(\vec{p}-\vec{q}) \delta^{s,r} |q, r\rangle$$

$$|p\rangle = \sqrt{2E_p} a_{\vec{p}}^{\dagger} |0\rangle$$

$$|q\rangle = \sqrt{2E_q} b_{\vec{q}}^{\dagger} |0\rangle$$

$$P = \int \frac{d^3p}{(2\pi)^3} \sum_s \vec{p} (a_{\vec{p}}^{\dagger} a_{\vec{p}}^s + b_{\vec{p}}^{\dagger} b_{\vec{p}}^s)$$

Observables: $O(x) = \bar{\psi}(x) \Gamma \psi(x)$

we want $[O(x), O(y)] = 0$ for $(x-y) < 0$